# Geometrical and Statistical Properties of Evolution Equations on Singular Spaces.

Andrei I. Shafarevich

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# Outline



2 Wave equation on the "hybrid space"

### 3 Time-dependent Schrödinger equation

- Gaussian beams
- Scattering on a manifold
- Statistics of the Gaussian packets
- Examples of exponential asymptotics

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# Singular spaces (decorated graphs)

Decorated graph  $\Gamma$  — topological space, obtained from a graph via replacing vertices by smooth manifolds  $M_k$ , dim $M_k \leq 3$ .



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Spectral theory:

L.D. Faddev, B.S. Pavlov, Theor. Math. Phys, 1983, 55(2), 257-268.

B.S. Pavlov, Theor. Math. Phys, 1984, 59(3), 345-353.

P.Exner, P. Sěba.J. Math. Phys. 28 (1987), 386-391.

J. Brüning and V.A. Geyler, M. Lobanov, S. Roganova, A. Tolchennikov, etc.

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## Laplacian

Edges  $\gamma_j$  of the graph are smooth parametrized curves,  $M_k$  are complete Riemannian manifolds.

Definition of the Laplace operator  $\Delta$ : 2 conditions.

- Δ is self-adjont;
- If *M* is a disconnected then



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Definition of the Laplace operator  $\Delta$ : 2 conditions.

- $\Delta$  is self-adjont;
- If *M* is a disconnected then

$$\Delta = \oplus_j \frac{d^2}{dz^2} \oplus_k \Delta_k.$$

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## Laplacian

Formal definition. Consider the direct sum

$$\Delta_0 = \oplus_j rac{d^2}{dz_j^2} \oplus_k \Delta_k$$

with the domain

$$H^2(\Gamma) = \oplus_j H^2(\gamma_j) \oplus_k H^2(M_k).$$

 $\Delta$  is a self-adjoint extension of the restriction  $\Delta_0|_L$ , where

$$L = \{\psi \in H^2(\Gamma), \quad \psi(q_s) = 0\}.$$

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# Coupling conditions

Boundary conditions: on the edges consider  $\psi(q_j)$ ,  $\psi'(q_j)$ . On the manifolds functions have singularities of the Green function type

$$\psi = a_j F(x) + b_j + o(1),$$

$$F = \begin{cases} -\frac{c}{2\pi} \log \rho, & \dim M = 2\\ \frac{c}{4\pi\rho}, & \dim M = 3. \end{cases}$$

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# Coupling conditions

Vector 
$$\xi = (u, v), u = (\psi'(q_1), \dots, \psi'(q_N), a_1, \dots, a_N),$$
  
 $v = (\psi(q_1), \dots, \psi(q_N), b_1, \dots, b_N).$   
In  $\mathbb{C}^{2N} \oplus \mathbb{C}^{2N}$  consider standard skew-Hermitian form

$$[\xi^1,\xi^2] = \sum_{j=1}^{2N} (u_j^1 \bar{v}_j^2 - v_j^1 \bar{u}_j^2).$$

and fix the Lagrangian plane L.

Coupling conditions

$$\xi \in L, \quad -i(E+U)u + (E-U)v = 0,$$

U is unitary matrix.

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Local coupling conditions — for each point separately:

$$L = \oplus_q L_q$$

For each point of gluing  $q \xi_q = (u_q, v_q)$ ,  $u_q = (\psi'(q), a)$ ,  $v_q = (\psi(q), b)$ ,

$$-i(E+U_q)u_q+(E-U_q)v_q=0,$$

 $U_q$  is a unitary 2 × 2-matrix.

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## Example: wave equation on the hybrid space

Example: wave equation on the hybrid space. The hybrid space  $\mathbb{R}^3$  with a half-line  $\mathbb{R}_+$ , coupled at the point z = 0, x = 0.



### Wave equation

#### Wave equation

$$\begin{split} \psi_{tt} &= \Delta \psi \\ \psi|_{t=0} &= \psi_0(z), \quad \psi_t|_{t=0} &= \psi_0'(z), \quad \psi_0 \in C_0^\infty(\mathbb{R}_+). \end{split}$$

Solution

$$\psi = \left\{egin{array}{ll} \psi_0(t+z) + \chi(t-z), & z\in\mathbb{R}_+, \ rac{f(t-|x|)}{|x|}, & x\in\mathbb{R}^3. \end{array}
ight.$$

Functions  $\chi(z)$ , f(z) can be expressed explicitly in terms of  $\psi_0$  and coupling matrix U.

## Examples

Complete reflection f = 0.

$$U = \begin{pmatrix} e^{i\varphi} & 0\\ 0 & e^{i\theta} \end{pmatrix}, \quad L = L_1 \oplus L_2.$$
$$\chi(z) = \frac{-i\frac{\partial}{\partial z}\cos\frac{\theta}{2} + i\sin\frac{\theta}{2}}{-i\frac{\partial}{\partial z}\cos\frac{\theta}{2} - i\sin\frac{\theta}{2}}\psi_0(z).$$

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# Examples

Complete transmission  $\chi = 0$ 

$$U=rac{1}{1+4\pi} egin{pmatrix} 1-4\pi & 4\sqrt{\pi}e^{iarphi}\ -4\sqrt{\pi}e^{-iarphi} & 1-4\pi \end{pmatrix}.$$
  $f(z)=rac{ie^{iarphi}}{2\sqrt{\pi}}\psi_0(z).$ 

Further examples (Anna V. Tsvetkova): explicit formulas for the wave equation on a space of constant curvature ( $S^2$ ,  $S^3$ ,  $H^2$ ,  $H^3$ ) coupled to the half-line.

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## Gaussian packet

Let  $\Gamma$  be a decorated graph. Cauchy problem for Schrödinger equation

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2}\Delta \psi + V\psi$$

$$\psi|_{t=0} = A_0 e^{\frac{iS_0(z)}{h}}.$$

$$S_0 = p_0(z - z_0) + \frac{1}{2}Q_0(z - z_0)^2, \quad \Im p_0 = 0, \quad \Im Q_0 > 0.$$

Asymptotics as  $h \rightarrow 0$ ; narrow Gaussian packet (squeezed state) located on the edge. Certain technical conditions on V, metrics and  $S_0$ .

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# Solution for small t

#### Solution for small *t*:

Assertion

$$\psi(z,t,h) = A(t)e^{\frac{iS(z,t)}{h}} + O(\sqrt{h}).$$

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$$S(z,t) = \sigma(t) + P(t)(z - Z(t)) + \frac{1}{2}Q(t)(z - Z(t))^{2}.$$

(P(t), Z(t)) is the solution of the classical Hamiltonian system

$$\dot{z} = \frac{\partial H}{\partial p}, \quad \dot{p} = -\frac{\partial H}{\partial z}, \quad H = \frac{1}{2}p^2 + V(z).$$

$$P(0) = p_0, \quad Z(0) = z_0.$$

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# Scattering on a manifold

Let  $\Gamma$  be a half-line, connected with the manifold M in a single point q.



#### Figure: Scattering on a manifold

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# Scattering on a manifold

#### Theorem

For certain time interval solution has the form

$$\psi = \begin{cases} A(t)e^{\frac{iS(z,t)}{h}}, & z \in \mathbb{R}_+, \\ K_{\Lambda_t}[B(x,t)], & x \in M. \end{cases} + O(\sqrt{h}) \end{cases}$$

K is the Maslov canonic operator on the isotropic manifold  $\Lambda_t$  with complex germ. A and B can be expressed explicitly in terms of the coupling matrix U.

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# Scattering on a manifold

#### Support of the solution on M.

At the instant of scattering  $t_0$  consider the sphere in  $T_q^*M$ :  $\Lambda_0: |p| = |P(t_0)|$ . Consider the flow  $g_t$  of the classical Hamiltonian system on M with the Hamiltonian  $H = \frac{1}{2}|p|^2 + V$ . Isotropic manifold  $\Lambda_t = g_t \Lambda_0$ ; as  $h \to 0$   $supp \psi \to \pi(\Lambda_t)$ ,  $\pi$  is the natural projection.

In general position  $\psi$  is localized near the surface of codimension 1 (Gaussian packet near the hypersurface). For V = 0 the support coincides with the geodesic sphere.

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# Scattering on a manifold



#### Figure: Support of the solution

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# Number of the Gaussian packets

Let  $\Gamma$  be a decorated graph with finite edges. For arbitrary finite t the solution has the form

$$\psi = \sum_{j} \psi_{j} + O(\sqrt{h}),$$

where  $\psi_j$  are the Gaussian packets.

Let N(t) be the number of packets, localized on the edges of  $\Gamma$  (not on the manifolds).

Let  $t_j$  be times of passage of the trajectories along the edges of the graph and between gluing points on the manifolds.

Assumption: there is a finite number of times  $t_1, \ldots, t_M$ .

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# Statistics of Gaussian packets



#### Figure: Times of connection

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# Statistics of Gaussian packets

#### Theorem

Let  $t_i$  be linearly independent over  $\mathbb{Q}$ . As  $t \to \infty$ 

$$N(t) = Ct^{M-1} + o(t^{M-1}).$$

### Constant C

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#### Theorem

For almost all t<sub>j</sub>

$$C = \frac{\sum_{edges} t_j}{2^{N-2}(M-1)! \prod_{all} t_j}.$$

Here N — number of points of gluing, M — number of times  $t_i$ .

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## Distribution of the packets

Uniform distribution. Let  $\delta$  be a segment on arbitrary edge of  $\Gamma$ .  $N_{\delta}$  is a number of packets, located on  $\delta$ .

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# Distribution of the packets

#### Uniform distribution.

Let  $\delta$  be a segment on arbitrary edge of  $\Gamma.$ 

 $N_{\delta}$  is a number of packets, located on  $\delta$ .

Theorem

$$\lim_{t\to\infty} \frac{N_{\delta}(t)}{N(t)} = \frac{t_{\delta}}{\sum_j t_j}.$$

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Ideas of the proofs.

1. Reduction to the problem of statistic of classical particles on metric graphs.

2. Application of the previous results (V. Chernyshev, A.Sh.) obtained for the case of metric graphs.

3. Reduction to the certain problem of analytic number theory - counting the number of integral points in a large simplex.

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# Hardy

Example: Star with three edges — Hardy - Littlewood formula (1927).

$$N(t) = \frac{t^2}{8} \frac{t_1 + t_2 + t_3}{t_1 t_2 t_3} + \frac{t}{2} (\frac{1}{t_1} + \frac{1}{t_2} + \frac{1}{t_3}) + o(t).$$



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# The cylinder with the segment

Examples of spaces with infinite number of times (V. Chernyshev,

A. Tolchennikov)



Figure: The segment glued to the cylinder in two points

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$$N(t) = C \exp\left(\sqrt{\frac{2}{3R}}\pi\sqrt{t}(1+o(1))
ight).$$

Here R is the radius of the cylinder.

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### Flat torus with the segment



#### Figure: The segment glued to the torus in two points

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$$N(t) = C \exp\left(3\sqrt[3]{rac{5\pi}{8ab}}\zeta(3)t^{2/3}(1+o(1))
ight).$$

Here a, b are lengths of the fundamental cycles of the torus. For 3D torus

$$N(t) = C \exp\left(4\sqrt[4]{rac{\pi}{3abc}}\zeta(4)t^{3/4}(1+o(1))
ight).$$

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Gaussian beams Scattering on a manifold Statistics of the Gaussian packets Examples of exponential asymptotics

The proof uses the following result.

Theorem (V.E. Nazaikinskii, 2013). Let N(t) be the number of non-negative integer solutions of inequality  $\sum_{i=1}^{\infty} \lambda_i N_i \leq t$  and for sequence  $\lambda_j$  a counting function  $\rho(\lambda) = \sharp(j|\lambda_j \leq \lambda)$  has asymptotics

$$\rho = c_0 \lambda^{1+\gamma} (1 + O(\lambda^{-\varepsilon})), \quad \varepsilon > 0.$$

Then

$$\log \mathsf{N}(t) = (\gamma+2) \left( \frac{c_0 \Gamma(\gamma+2) \zeta(\gamma+2)}{(\gamma+1)^{\gamma+1}} \right)^{\frac{1}{\gamma+2}} t^{\frac{\gamma+1}{\gamma+3}} (1+o(1)).$$

as t goes to infinity

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